FLOW DISTURBANCES DUE TO WALL IRREGULARITIES IN TWO-DIMENSIONAL SUPERSONIC FLOW

The Magnitude and Lateral Extent of Disturbance For Mach Numbers Between 1.05 and 1.25

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FOREWORD

This report was prepared under Expenditure Order 903-1128 by Captain Jean G. Goppert, as project engineer in the Wind Tunnel Branch, at the request of the Chief of the Project Unit, Wind Tunnel Branch.

Included among those who cooperated in reviewing and editing the report were Dr. B. H. Goethert and Mr. R. R. Kassner of the Wind Tunnel Branch.

ABSTRACT

A method is presented for computing the lateral extent of a weak shock wave due to a small wedge-shaped irregularity of a wall in two-dimensional supersonic flow. Computations have been carried out over the Mach number range 1.05 to 1.25 for angular irregularities of the wall of 0.1 to 0.4 degrees. The entropy change in the shock wave as well as the influence of the wall boundary layer were assumed to be negligible. Computations were carried out for the straight part of the shock wave and for such points along the curved part of the shock wave at which the local Mach number change, Δ M, across the shock wave, is reduced to values of 0.001, 0.002, 0.003, 0.004, and 0.005.

PUBLICATION REVIEW

Manuscript Copy of this report has been reviewed and found satisfactory for publication.

FOR THE COMMANDING GENERAL:

ACK A. GIBBS

doldnel, USAF

Chief, Aircraft Laboratory

Engineering Division

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SYMBOLS

- M Mach number
- My Free stream Mach number
- M_2 $(M_1-\Delta M_2)$ Mach number behind the shock wave, but in front of the expansion fan
- M_3 $(M_1-\Delta M_3)$ Mach number along some expansion wave at which the flow has attained a Mach number arbitrarily close to M_1 .
- Length of the wedge
- 8 Wedge angle
- Δh Maximum wedge thickness
- $oldsymbol{ heta}$ Angle measured counterclockwise from the horizontal
- θ_1 Mach angle of flow at M₁
- $heta_2$ Angle of first expansion wave
- $heta_3$ Angle of the expansion wave along which the flow has attained the Mach number $exttt{M}_3$
- € Angle of the shock wave at any point
- ϵ_2 Angle of the shock wave before intersection with the expansion fan
- Mach angle measured from the local flow direction
- ν_{l} Deflection angle required to change the flow Mach number from M = 1.0 to M_{1} in a Meyer expansion
- ν Deflection angle required to change the flow Mach number from M = 1.0 to M_2
- r Radius from the rear corner of the wedge
- r₂ Radius to the point of intersection of the shock wave and the first expansion wave
- Radius to the point on the shock wave where the Mach number behind the shock has attained the value Mz

- y Lateral distance (perpendicular to free stream)
- Distance from wall to point of intersection of shock wave and first expansion wave
- Distance from wall to point along shock wave where the Mach number behind the shock has attained the value M_3
- $\gamma = \frac{C_p}{C_w}$ Ratio of specific heats

SECTION I

PURPOSE

The purpose of this investigation is to formulate an estimate within practical accuracy of the magnitude and extent of the disturbance due to a small irregularity of a wall in two dimensional supersonic flow. Consideration is limited to low supersonic Mach numbers and small irregularities, for which only attached shock waves appear and the shock losses are negligible. Boundary layer influences are not included, but it can be expected that they will weaken the intensity of the flow disturbance.

SECTION II

DERIVATION

The problem has been simplified by considering the special case of a plane surface with a wedge shaped irregularity in supersonic flow.

Figure I shows the wedge shaped irregularity AB. Upstream of A and downstream of B the wall is assumed to be plane and parallel to the undisturbed flow. A shock wave originates at A, changing the flow angle from 0 to 8 and the Mach number from M_1 to M_2 . These flow conditions, however, exist only within the triangle ABP2. At point P2 (wall distance Y2) the shock wave is intersected by the first of the expansion waves which start from the corner at B. Beyond P2 the shock wave is curved due to intersection with the expansion waves and the intensity of the shock wave is reduced. Therefore, the Mach number change, ΔM_2 , $(M_1 - M_3)$, across the shock wave, occurring at a point, P3 (wall distance Y3), on the curved part, is smaller than the change ΔM_2 across the straight part of the shock.

The object of the following computation is to determine:

- a. Mach number disturbance, Δ Mz, across the curved part of the shock as a function of distance Yz.
- b. Mach number disturbance across the straight part (equal to disturbance at wall) and distance Y_2 .

Curved Part of the Shock wave

Figure 2 shows an element of the shock wave between two expansion waves having the angles θ and θ -d θ with the free stream direction. The shock wave element forms the angle ϵ with the free stream direction.

Therefore, since $(\theta - \epsilon)$ is the angle between the shock wave element and an expansion wave, and r is the radius from the corner, B, we have:

$$\frac{\mathrm{d}\mathbf{r}}{-\mathrm{rd}\boldsymbol{\theta}} = \cot \left(\boldsymbol{\theta}_2 - \boldsymbol{\epsilon}\right) \tag{1}$$

integrating from r2 to r3 (See Figure 1)

$$\ln \frac{r_3}{r_2} = \int_{\theta_3}^{\theta_2} \cot (\theta - \epsilon) d\theta \qquad (2)$$

Thus the first problem is to set up $(\theta - \epsilon)$ in terms of θ . For a given initial Mach number, M₁, and for assumed values of the shock wave angle, ϵ , the Mach number, M, behind the shock wave can be computed from the equation (Eq. 4.33, Ref. 3).

$$M^{2} = \frac{1 + \frac{\gamma - 1}{2} M_{1}^{2}}{\gamma M_{1}^{2} \sin^{2} \epsilon - \frac{\gamma - 1}{2}} + \frac{M_{1}^{2} \cos^{2} \epsilon}{1 + \frac{\gamma - 1}{2} M_{1}^{2} \sin^{2} \epsilon}$$
(3)

The proper range of assumed values of ϵ was determined from a plot of δ vs ϵ for constant values of M_1 , taken from Table I of Ref. 2. By neglecting the entropy changes in the weak oblique shocks, these computed values of M_1 , together with Table II of Ref. 1, for issentropic flow, can be used to obtain θ , which is the angle of the expansion wave along which the Mach number is M_1 . That is:

$$\theta = \alpha + (\nu_1 - \nu) \qquad (4)$$

where α is the Mach angle corresponding to the Mach number, M, and $(\nu_l - \nu)$ is the flow inclination along the Mach wave with respect to the direction of the undisturbed flow with M = M₁. (See Figure 3.)

Thus we have calculated values of θ and ϵ corresponding to various values of M. It is now possible to compute and plot $(\theta^-\epsilon)$ vs θ^- .

If the plot of $(\theta - \epsilon)$ vs θ can be approximated by a straight line over the desired range of θ , then the integral (Eq. 2) can be evaluated very easily since we can substitute

$$\theta - \epsilon = \frac{\theta + c}{a} \tag{5}$$

where the constants a and c can be evaluated from the geometry of the line as

$$a = \frac{d\theta}{d(\theta - \epsilon)}$$

$$C = -\theta_1$$

Similarly, a slight curvature of the plot $(\theta - \epsilon)$ vs θ can be approximated by two straight lines and integrated in two intervals. If, however, the curvature is too great, then a graphical integration seems to be the best procedure. Herely plot a curve of cot $(\theta - \epsilon)$ vs θ , compute the area under the curve between the limits θ_2 and θ_3 , and take the antilog to get $\frac{r_3}{r_2}$.

The angle θ_2 is defined as the angle of the first expansion wave of the Prandtl-Meyer expansion, (Figures 1 and 3).

Thus
$$\theta_2 = \alpha_2 + \delta$$
 (6)

 $heta_3$ is defined as the angle of the expansion wave along which the flow has reached any arbitrary Mach number, M₃, which is determined by the choice of a value of Δ M₃ across the shock wave. After a selection of Δ M₃ is made, θ_3 can be computed by means of Eq. 3 and 4.

Straight Part of the Shock Wave

The next step is to compute the value of $\frac{r_2}{1}$ which is the radius from

the expansion corner, B, to the point of the intersection of the shock wave with the first expansion wave, expressed in terms of the length of the wedge. This is a simple problem in trigonometry since the initial shock wave angle, ϵ_2 , and the Mach number behind the shock, M, can be determined from tables in Ref. 2, Eq. 3. Since the direction of a Mach wave is given by the Mach angle, which depends upon the known Mach number of the flow, we can, by referring to Figure 1, see that the law of sines will give the desired radius, ϵ_2

(1) $\frac{\mathbf{r}_2}{\mathbf{I}} = \frac{1}{\cos \delta} \cdot \frac{\sin (\epsilon_2 - \delta)}{\sin (\theta_2 - \epsilon_2)}$ (7)

(2)
$$\theta_2 - \epsilon_2 = \alpha_2 - \epsilon_2 + \delta$$

(3)
$$\frac{\mathbf{r}_2}{\mathbf{1}} = \frac{\sin(\epsilon_2 - \delta)}{\cos \delta \sin(\alpha_2 - \epsilon_2 + \delta)}$$

It is now a simple step to determine $\frac{Y_3}{2}$ (See Figure 1)

$$\frac{Y_3}{1} = \frac{r_3}{1} \sin \theta_3$$

T:

RESULTS AND DISCUSSION

Computations were carried out for the Mach number range 1.05 to 1.25, and for wedge angles of 0.1, 0.2, 0.3, and 0.4 degrees, using arbitrary values of ΔM_3 = .001 to .005. The results are presented in Figures 4 and 5.

The probable error in Y_2 and Y_3 may be as great as 5% near M_1 = 1.05, with considerable improvement as the Mach number increases. This large error is due to the extremely small difference between the shock wave angle and the angle of the expansion waves. However, it must be remembered that only the approximate order of magnitude of the lateral extent of the disturbance is all that was to be obtained from this calculation.

The case of negative values of the wedge angle, δ , was checked at M_1 = 1.15. The values of Y_2 were found to be of the same order of magnitude but slightly smaller than the values for positive angles. The difference was so small, however, that further computations did not appear to be warranted.

The results have been presented in Figure 5 as a single family of curves of Δ M_Z vs Y_Z/ $\mathbf{1}$ δ for lines of constant Mach number, M₁. The computations showed that Y_Z/ $\mathbf{1}$ was proportional to δ within the investigated range of wedge angles and Mach numbers and within the accuracy of these computations. It should be noted, also, that the values given in Figure 5 are not valid for values of Y_Z < Y_Z, (see Figure 1), since the disturbance Δ M is constant from Y = 0 to Y = Y_Z.

The distance, Y_2 , as well as the disturbance ΔM_2 , valid for $0 \le Y \le Y_2$ has been plotted in Figure 4 with the wedge angle δ as a parameter. For the purpose of interpolating between the curves, it may be remarked that at constant Mach number, M_1 , ΔM_2 is approximately proportional to δ , while $Y_2/\Delta h$ is approximately proportional to $1/\delta^2$. Furthermore, for a given wedge, the distance Y_2 varies approximately with the Mach number according to M_1^2 -1.

APPLICATION OF RESULTS

There are two problems which can be solved quite easily using Figure 4.

a. Problem 1. Given a wall configuration, δ and $\mathbf{1}$, and a free stream Mach number, M_1 . Find the strength of the disturbance, Δ M_3 , at some given distance, Y_3 , from the wall.

Solution. Compute $\frac{Y_3}{1 \cdot 8}$, and read the value of Δ M₃ from the curve on Figure 5 corresponding to the proper value of M₁.

b. Problem 2. Given the maximum allowable disturbance, Δ M₃, at a distance, Y₃, from the wall, and a free stream Mach number, M₁. Find the allowable height of wall disturbance Δ h = £8 which would satisfy the maximum allowable disturbance requirements.

Solution. From Figure 5, find the value of $\frac{Y_3}{1\cdot 8}$ corresponding to the given values of M₁ and Δ M₃. Then, since Y₃ is given, this determines the allowable disturbance height, 1 · 8.

If, in the problem to be solved, the wall has a curved irregularity instead of the assumed abrupt wedge-shaped irregularity, the results given in this report will be conservative since the initial shock wave will be split up into a number of compression waves which will probably converge into a shock wave, but the intensity will be less than for the wedge.

The same effect could be expected in considering the boundary layer, since the boundary layer would tend to round off the corners of the wedge.

SUMMARY

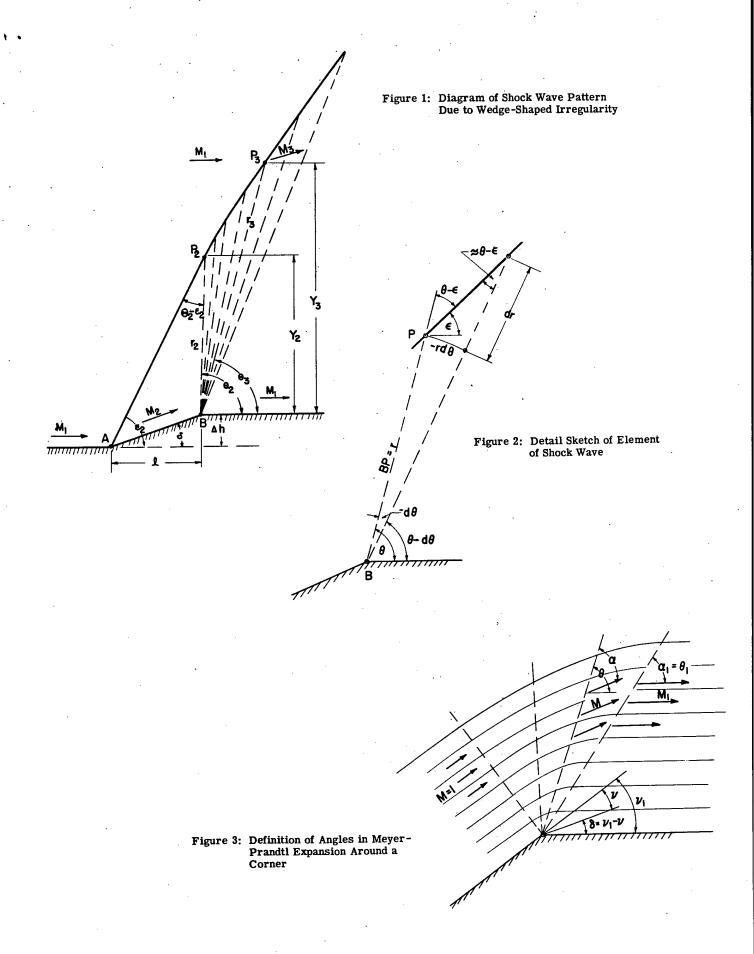
The local Mach number disturbances in two dimensional flow, due to a wedge shaped wall irregularity, are calculated for various lateral distances from the wall for the Mach number range between M = 1.05 and 1.25. (See Figures 4 and 5.)

When the height of the disturbance $\Delta h = \delta \cdot \mathbf{1}$ is kept constant, the lateral extent of the region between the initial shock wave and the first expansion wave increases inversely with the square of the wedge angle of the wall irregularity and varies with the Mach number according to $\mathbb{M}_1^2 - 1$. (See Figure 4.)

In the region behind the first intersection of initial shock wave and expansion wave, the local Mach number disturbance is nearly independent of Mach number and wedge angle, but rather depends mainly on the ratio of the lateral distances to the disturbance height. (See Figure 5.)

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- 2. NACA TN 1673 Tables and Charts of Flow Parameters Across Oblique Shocks, by Mary M. Neice, August 1948
- 3. Liepmann & Puckett "Aerodynamics of a Compressible Fluid"
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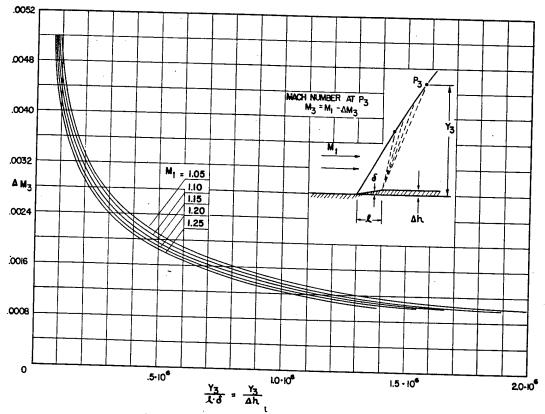


Figure 4: Lateral distance to first intersection of shock wave and expansion fan and Mach number disturbance at the wall for small wedge-shaped irregularity of a two-dimensional wall.

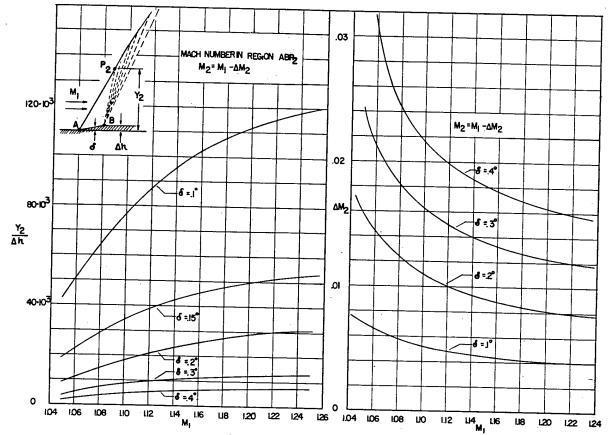


Figure 5: Lateral variation of Mach number disturbance due to a small wedge-shaped irregularity of a two-dimensional wall. Region downstream of first intersection between shock wave and expansion wave.